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A "grazing particle" just touches the crystal during simultaneous betatron & synchrotron oscillation extrema.

$$a_x + |\eta| a_s = |x_c|$$



### The ONLY page of equations



Thus the grazing angle depends linearly on the synchrotron amplitude  $a_s$  according to

$$x'_G = -\frac{\alpha}{\beta} x_c + \operatorname{sgn}(x_c) \operatorname{sgn}(\eta) g a_s$$
(15)

where the linear slope of grazing angle with respect to synchrotron amplitude is

$$\frac{dx'_G}{da_s} = \operatorname{sgn}(x_c)\operatorname{sgn}(\eta)g \tag{16}$$

The grazing function g that enters these equations is an optical quantity defined as

$$g \equiv \left(\frac{\alpha}{\beta}\eta + \eta'\right) \tag{17}$$

The grazing function is thus revealed to be just

$$g = \sqrt{\beta} \, \eta'_N$$

and the rigorous general synchrobet atron condition  $g=0\ {\rm is}\ {\rm just}$ 

$$\eta'_N = 0$$



### More realistic condition on g



## Efficient collimation: the grazing angle must be within the crystal acceptance angle for "all" synchrotron amplitudes.

$$|g| < \frac{\sigma'_A}{n_{max} \left(\sigma_p / p\right)}$$

$$\sigma'_A$$
 [µrad] ~ 2 channeling at 7 TeV

$$\sim$$
 10 channeling at 0.1 TeV

$$\sim 100$$
 volume reflection at any energy

	α	eta [m]	$\eta$ [m]	$\frac{\eta'}{[10^{-3}]}$	$g$ $[10^{-3}]$	E [TeV]	$\frac{\sigma_p/p}{[10^{-3}]}$	$\sigma'_G$ [ $\mu$ rad]
RHIC SPS Tevatron LHC	-26.5 -2.21 -0.425 1.94	$1155.0 \\96.1 \\67.5 \\137.6$	$-0.864 \\ -0.880 \\ 1.925 \\ 0.559$	$-16.2 \\ -19.0 \\ 15.0 \\ -8.9$	$3.6 \\ 1.2 \\ 2.9 \\ -1.0$	$\begin{array}{c} 0.10 \\ 0.12 \\ 0.98 \\ 0.45 \\ 7.0 \end{array}$	$\begin{array}{c} 0.50 \\ 0.40 \\ 0.14 \\ 0.31 \\ 0.11 \end{array}$	$1.81 \\ 0.48 \\ 0.41 \\ 0.31 \\ 0.11$



UA9 looks ok: the grazing angle changes by only 0.48  $\mu$ rad per  $\sigma_{p}/p$ , despite large negative dispersion at the crystal.



Large  $\sigma_p/p$ : grazing angle changes by 5.4 µrad over 3  $\sigma_p/p$ . Helps explain disappointing performance?

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#### **Tevatron T980**





Flattened slopes: grazing angle changes by 0.41  $\mu$ rad per 3  $\sigma_p/p$ . Measured optics values also available.

LARP









Small  $\sigma_p/p$ , small channeling acceptance: grazing angle changes by only 0.11 (0.31) µrad per  $\sigma_p/p$  at 7 (0.45) TeV.

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## How does *g* propagate?

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A thin dipole: 
$$\Delta g = \left(\Delta \eta' - \frac{\eta \,\Delta b'}{b}\right) = \Delta \theta$$
  
A thin quad:  $\Delta g = \eta \left(\frac{\Delta \eta'}{\eta} - \frac{\Delta b'}{b}\right) = 0$ 





<u>A matched FODO cell</u> short dipoles thick quads:

<u>A matched FODO cell</u> long dipoles thick quads:

g is small when the normalized dispersion is almost constant

$$g = \sqrt{\beta} \eta'_N$$

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Simple prediction vs numerical testing



# Predict $|g|_{max} \approx \theta \frac{3}{4\sqrt{2}} \frac{\sqrt{1+C^2}}{S^2 C} \left( (2-S)\sqrt{1+S} - (2+S)\sqrt{1-S} \right)$

where  $S \equiv \sin \phi/2$ ,  $C \equiv \cos \phi/2$ 

for a matched FODO cell, phase advance per cell  $\,\phi,\,half$  cell length L, half cell bend angle  $\,\theta$ 

#### In reasonable agreement with numerical testing:

$$g_{max} \approx 0.427 \ L^0 \ \theta^1 \tag{51}$$

when the phase advance per full-cell is 90 degrees. There is no dependence on the cell length! A fair rule of thumb is that

$$g_{max} \approx \theta/2$$
 (52)

#### **CAUTION:** these best case results assume matched optics.

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## Conclusions



1) The grazing function g parameterizes the rate of change of total angle with synchrotron amplitude.

2) A pure optics function, it is related to the slope of the normalized dispersion. Ideal crystal value g = 0

3) g should be small enough that "all" synchrotron amplitudes are within the crystal acceptance angle.

4) This appears to be reasonable to achieve in practice, especially in VR mode, and at lower energies.

5) Place crystal away from betatron & dispersion waves, since they may increase g by an order of magnitude.

6) Planning for future crystals should include a grazing function analysis, both in design and in error analysis.