

# Adding RF-Multipoles to Mad-X

AL, LD, RdM, MG

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# Rationale

- Add RF-Multipolar fields to Mad-X
  - RF-Multipole: RF + Multipole
  - Properties of an RF-CAVITY
  - Properties of a multipole
- Ultimate objective: implement crab cavities with high-order RF multipolar errors
- Staged approach (conservative)

# RF-Multipole

is a Multipole with coefficients that change with  $z$

$$B_y(x, y, z) + iB_x(x, y, z) = \sum_{n=1}^N (B_n(r_0, z) + iA_n(r_0, z)) \cdot \frac{(x + iy)^{n-1}}{r_0^{n-1}}$$

$r_0$  : reference radius

$x, y, z$  : particle coordinates w.r.t. reference particle

Where  $B_n$  and  $A_n$  are the real part of a complex vector (phasor) rotating with angular velocity  $\omega_{\text{RF}}$

$$A_n(r_0, z) = \text{Re} \left\{ A_n(r_0) e^{j(k_{\text{RF}}z + \vartheta_n)} \right\} : \text{skew multipole component}$$

$$B_n(r_0, z) = \text{Re} \left\{ B_n(r_0) e^{j(k_{\text{RF}}z + \varphi_n)} \right\} : \text{normal multipole component}$$

$z = ct$  : longitudinal coordinate w.r.t. reference particle

$$\boxed{k_{\text{RF}}} = \frac{\omega_{\text{RF}}}{c} = \frac{2\pi}{\lambda_{\text{RF}}} : \text{RF wave number}$$

$\boxed{B_n(r_0)}$  : modulus of the phasor  $\vec{B}_n$

$\boxed{\vartheta_n}$  : phase of the phasor  $\vec{B}_n$

$\boxed{A_n(r_0)}$  : modulus of the phasor  $\vec{A}_n$

$\boxed{\varphi_n}$  : phase of the phasor  $\vec{A}_n$

# Two steps for the implementation

- 1) Add a **new element**: RF-Multipole
  
- 2) Modify **EFCOMP** to carry the additional information required by the RF-Multipole:
  - Frequency of the fundamental mode
  - Phase of each multipolar component

**Remark:** PTC implements a simplified version of the RF-Multipole (1 phase for all orders) that is not usable from Mad-X

# (1) New element

```
RFMULTIPOLE,  
  L=real, VOLT=real, LAG=real, HARMON=integer,  
  FREQ=real, LRAD=real, TILT=real,  
  KNL:={..., ..., ...}, ! Normal coefficients  
  KSL:={..., ..., ...}, ! Skew coefficients  
  PNL:={..., ..., ...}, ! Normal phases [2pi]  
  PSL:={..., ..., ...}; ! Skew phases [2pi]
```

RF-Multipole: RF + Multipole

- Properties of an RF-CAVITY
- Properties of a multipole

## (2) Extending EFCOMP

EFCOMP associates field and phase errors to an element

```
EFCOMP, ORDER:=integer, RADIUS:=real,  
  DKN:={dkn(0),dkn(1),dkn(2), ...},  
  DKS:={dks(0),dks(1),dks(2), ...},  
  DKNR:={dknr(0),dknr(1),dknr(2), ...},  
  DKSR:={dksr(0),dksr(1),dksr(2), ...},  
  FREQ:=real,  
  DPN:={dpn(0),dpn(1),dpn(2), ...},  
  DPS:={dps(0),dps(1),dps(2), ...};
```

Phase errors are always absolute

# Impact on the code (structure)

- EFCOMP:
  - Must be changed to take into account the new input vectors with phases, and the frequency of the fundamental mode
- Add a new element: RF-Multipole
  - Modify several modules: dict, node, ...
- Export to SixTrack
  - Modify module: 6track

Once these points achieved, we can start exporting to SixTrack

# Impact on the code (physics)

- Add a new element: RF-Multipole
  - Modify several modules: twiss, tracking, makethin, ...
  - Check closed orbit calculation (4d/6d?)
- Must check out the code already existing for CRABCAVITY :
  - It is implemented in tracking
  - It is not implemented in twiss
  - It is exported to SixTrack
  - What to do: We see two options:
    - Leave it as it is
    - Convert it to a RF-Multipole if needed



# Export to SixTrack

- We will adapt our export format to suit any SixTrack implementation of the RF-Multipole
- Preliminary discussions with RT, JB and FS indicate that a staged approach will be chosen (mainly because of limited resources):
  - RFQUAD, RFSEXTUPOLE, RFOCTUPOLE
  - Followed by RFMULTIPOLE at a later moment

# Conclusions

Proposed steps:

- Add a new element RF-Multipole
- Modify EFCOMP
- Develop the interface toward SixTrack.
- Twiss and tracking implementations could happen at a later time

# Appendix - Multipole Coefficients

## Multipole expansion

$$B_y(x, y) + iB_x(x, y) = \sum_{n=1}^N (B_n(r_0) + iA_n(r_0)) \cdot \frac{(x + iy)^{n-1}}{r_0^{n-1}} \quad [\text{T}]$$

$r_0$  : reference radius [m]

$x, y$  : particle coordinates w.r.t. reference particle [m]

$B_n(r_0)$  : normal multipole component at  $r_0$  [T]

$A_n(r_0)$  : skew multipole component at  $r_0$  [T]

## Mad-X integrated strength

$$\text{KNL}[\mathbf{n}] := l \cdot \frac{1}{B\rho} \frac{\partial^n B_y}{\partial x^n} = l \cdot \frac{1}{B\rho} \frac{B_{n+1}(r_0)}{r_0^n} n!$$

$$\text{KSL}[\mathbf{n}] := l \cdot \frac{1}{B\rho} \frac{\partial^n B_x}{\partial x^n} = l \cdot \frac{1}{B\rho} \frac{A_{n+1}(r_0)}{r_0^n} n!$$

$\text{KNL}[\mathbf{n}]$  : integrated nominal multipole strength of order  $n$  [ $\text{m}^{-n}$ ]

$\text{KSL}[\mathbf{n}]$  : integrated skew multipole strength of order  $n$  [ $\text{m}^{-n}$ ]

$l$  : length of the element [m]

$B\rho$  : magnetic rigidity [T m]

# Appendix - Field Errors

Absolute errors in Mad-X

$$\text{DKN} \equiv \text{KNL} \quad [\text{m}^{-n}]$$

$$\text{DKS} \equiv \text{KSL} \quad [\text{m}^{-n}]$$

Relative error for the multipole strength of order  $n$  in Mad-X

$$\text{KNL}_{[n]} = \text{DKNR}_{[n]} \cdot \text{KNL}_{[\text{ORDER}]} \cdot r_0^{(n-N)}$$

$$\text{KSL}_{[n]} = \text{DKSR}_{[n]} \cdot \text{KSL}_{[\text{ORDER}]} \cdot r_0^{(n-N)}$$

$\text{DKNLR}_{[n]}$  : relative error for normal component of order  $n$  [real]

$\text{DKSLR}_{[n]}$  : relative error for skew component of order  $n$  [real]

Where the relative errors are a scaling factor for  $b_n$  (and  $a_n$ ), relative to  $B_N(r_0)$

$$B_n(r) = B_N(r_0) b_n(r)$$

$$b_n(r) = b_n(r_0) \left( \frac{r}{r_0} \right)^{n-N}$$

$r_0$  : reference RADIUS [m]

$N$  : ORDER [integer]

$r$  : arbitrary radius [m]

$B_N(r_0)$  : main field at radius  $r_0$  [T]

$B_n(r)$  : normal multipole component at radius  $r$  [T]

$b_n(r)$  : normal relative multipole coefficient related to the main field  $B_N(r_0)$  [real]

# Appendix - Multipole Expansion

Given

$$\tilde{\mathbf{E}} = \mathbf{E}(x, y, z) e^{-j(\omega t + \phi)},$$

$$\tilde{\mathbf{H}} = \mathbf{H}(x, y, z) e^{-j(\omega t + \phi)}$$

- The magnetic strength [A/m] is converted into magnetic induction,  $B$  [T]

$$\mathbf{B} = \mu_0 \mathbf{H}$$

- The electric field and the magnetic strength are added, using the Lorentz force

$$\mathbf{E}_{\text{equiv.}} = \mathbf{E} + \mathbf{v} \wedge \mathbf{B}$$

that is,

$$\mathbf{E}_{\text{equiv.}} = \mathbf{E} + \det \begin{pmatrix} \hat{\mathbf{e}}_x & \hat{\mathbf{e}}_y & \hat{\mathbf{e}}_z \\ 0 & 0 & c \\ B_x & B_y & B_z \end{pmatrix} = \begin{pmatrix} E_x - cB_y \\ E_y + cB_x \\ E_z \end{pmatrix},$$

where  $\mathbf{v} = (0, 0, c)$ . Note that  $B_z = 0$  if the excited mode is TM.

Then the resulting field  $E$  can be expanded in either ways:

$$E_y + iE_x = \sum_{n=1}^N C_n (x + iy)^{n-1}$$
$$E_z = \sum_{n=0}^N D_n (x + iy)^n$$