

# Potential for Stochastic Cooling of Heavy Ions in the LHC

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# Outline

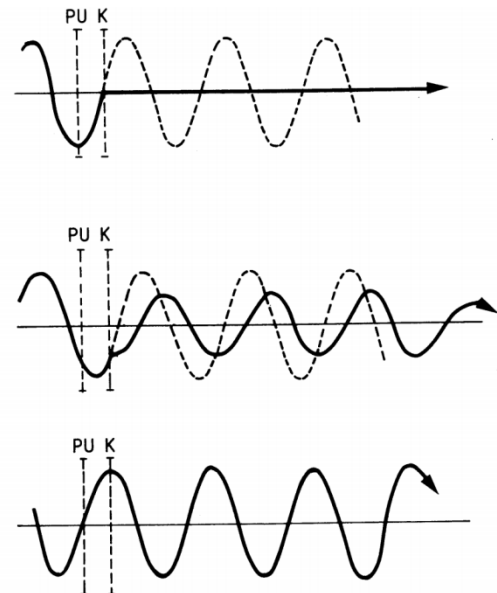
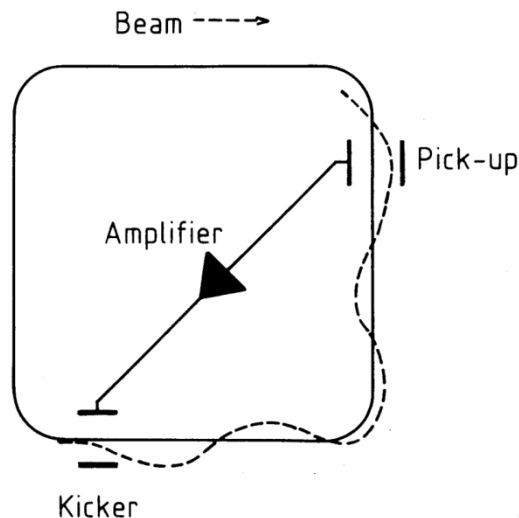
- Brief Introduction to Stochastic Cooling
- System and Performance at RHIC
- Data Analysis and Simulations from 2011 and 2013
  - Bunch-to-Bunch Differences
  - Beam Evolution and Tracking Simulations
- First Studies for a Stochastic Cooling System at LHC

# Stochastic Cooling – Basic Principle (1/5)

## Cooling aims to reduce the size and energy spread of a circulating particle beam

### Test particle picture:

- Particle performs betatron oscillations, due to position and angle errors.
- Cooling system is designed to damp these.
- Pick-up measures transverse position at each turn.
- Kicker applies angle correction proportional to position error at pick-up.
  - Synchronism between particle & signal!
  - Phase of pick-up and kicker is chosen to be  $(\lambda/4 + n \lambda/2)$ .
  - Signal has to take a short cut.



Oscillation completely cancelled.

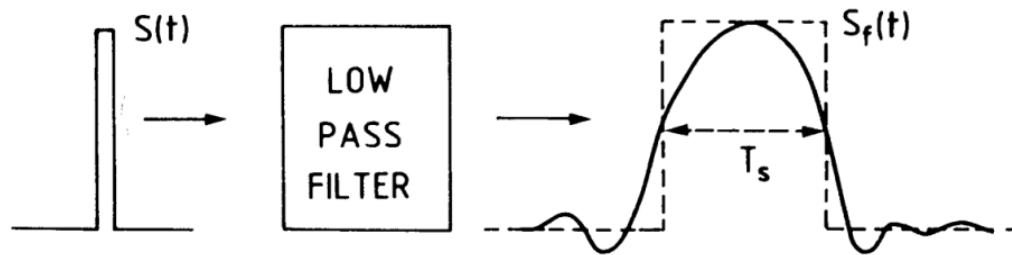
Oscillation partly cancelled.

Particle not affected.

# Stochastic Cooling – Basic Principle (2/5)

## Beam samples:

- Off-axis particle passing through pick-up induces short pulse.
- Finite bandwidth ( $W$ ) of the cooling system.  
→ Kicker signal broadened into pulse of length  $T_s = 1/(2W)$ .



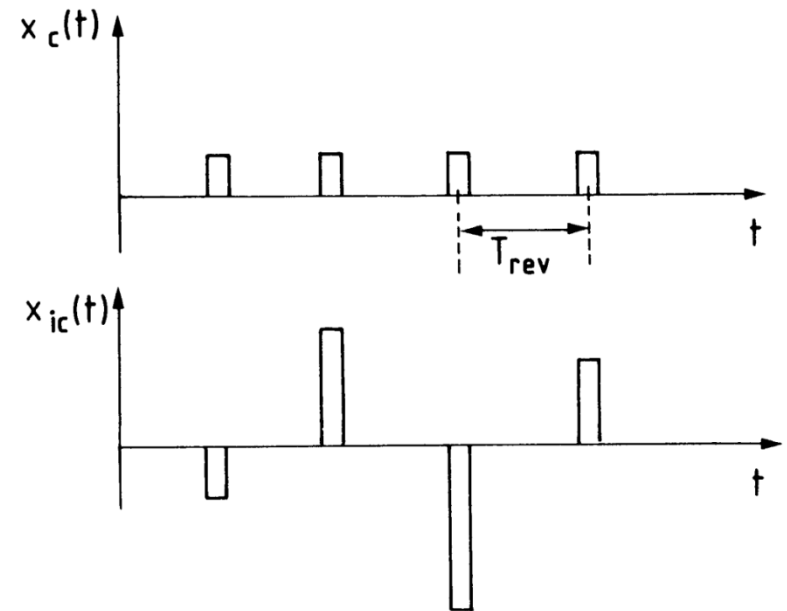
- Test particle passing system at  $t_0$  will be affected by kicks of all particles passing during the time interval  $t_0 \pm T_s / 2$ .  
→ Those particles belong to the **sample of the test particle**.

# Stochastic Cooling – Basic Principle (3/5)

## Test particle picture:

- $x$  is the error of the test particle.
- Corresponding correction at the kicker:  $\Delta x = -\lambda x$   
→ Corrected error:  
$$x_c = x - \lambda x$$
- Kicks of other sample members have to be added!

Coherent (systematic) signal of the test particle itself.

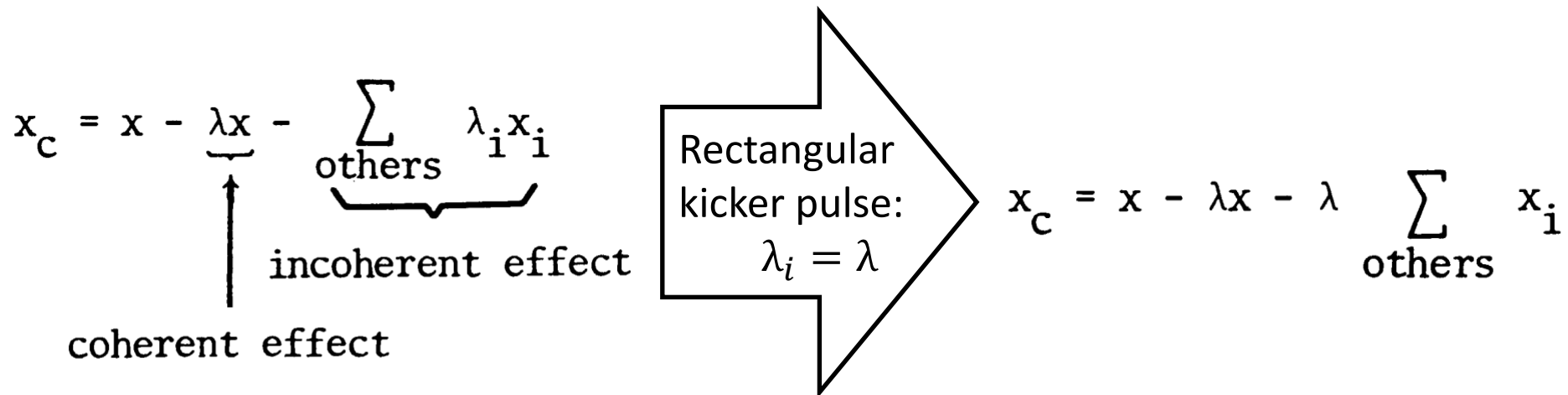


Incoherent (random) signal of the other sample particles.

→ Heating!

# Stochastic Cooling – Basic Principle (4/5)

- Only test particle present:
  - Correction at the kicker:  $\Delta x = -\lambda x$
  - Corrected error:  $x_c = x - \lambda x$
- Kicks of other sample members have to be added.



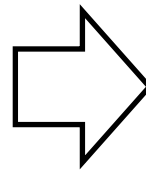
# Stochastic Cooling – Basic Principle (5/5)

- Rewrite the sum to include the test particle.
- Average sample error (the samples centre of gravity):  
 $N_s$  is number of particles per sample.

$$\langle x \rangle_s = \frac{1}{N_s} \sum_{\text{sample}} x_i$$

Test particle picture

$$x_c = x - \lambda x - \lambda \sum_{\text{others}} x_i$$



Sampling picture

$$x_c = x - (N_s \lambda) \langle x \rangle_s$$

**Cooling System measures average sample error and applies a correcting kick ( $\propto \langle x \rangle_s$ ) to the test particle!**

# Crude Cooling Rate Approximation (1/2)

$$\Delta x = -(\lambda N_s) \langle x \rangle_s \equiv -g \langle x \rangle_s$$

- $g$  can be interpreted as the fractional correction per turn  $\rightarrow g \leq 1$ .
- $g = \lambda N_s = 1$  is an estimate for the upper limit.
- Assume incoherent effect averages to zero:

$$\Rightarrow \Delta x = -\frac{1}{N_s} x$$

Assume exponential damping with the number of turns  $n$ :

$$x = x_0 \exp(-n / \tau_n)$$

Cooling rate per turn.  $\Rightarrow \frac{1}{\tau_n} = -\frac{1}{x} \frac{dx}{dn} \cong -\frac{1}{x} \frac{\Delta x}{\Delta n} = -\frac{\Delta x}{x} = \frac{1}{N_s}$

↑  
since  $\Delta n = 1$  turn



# Crude Cooling Rate Approximation (2/2)

$$\text{Cooling rate per turn: } \frac{1}{\tau_n} = \frac{1}{N_s}$$

For a coasting beam:

$$\text{Number of samples per beam: } n_s = T_{rev} / T_s$$

$$\text{Number of particles per sample: } N_s = N / n_s = NT_s / T_{rev} = N / (2WT_{rev})$$

$$\text{Cooling rate per second: } \frac{1}{\tau} = \frac{1}{T_{rev}} \frac{1}{\tau_n} = \frac{2W}{N}$$

For a bunched beam:

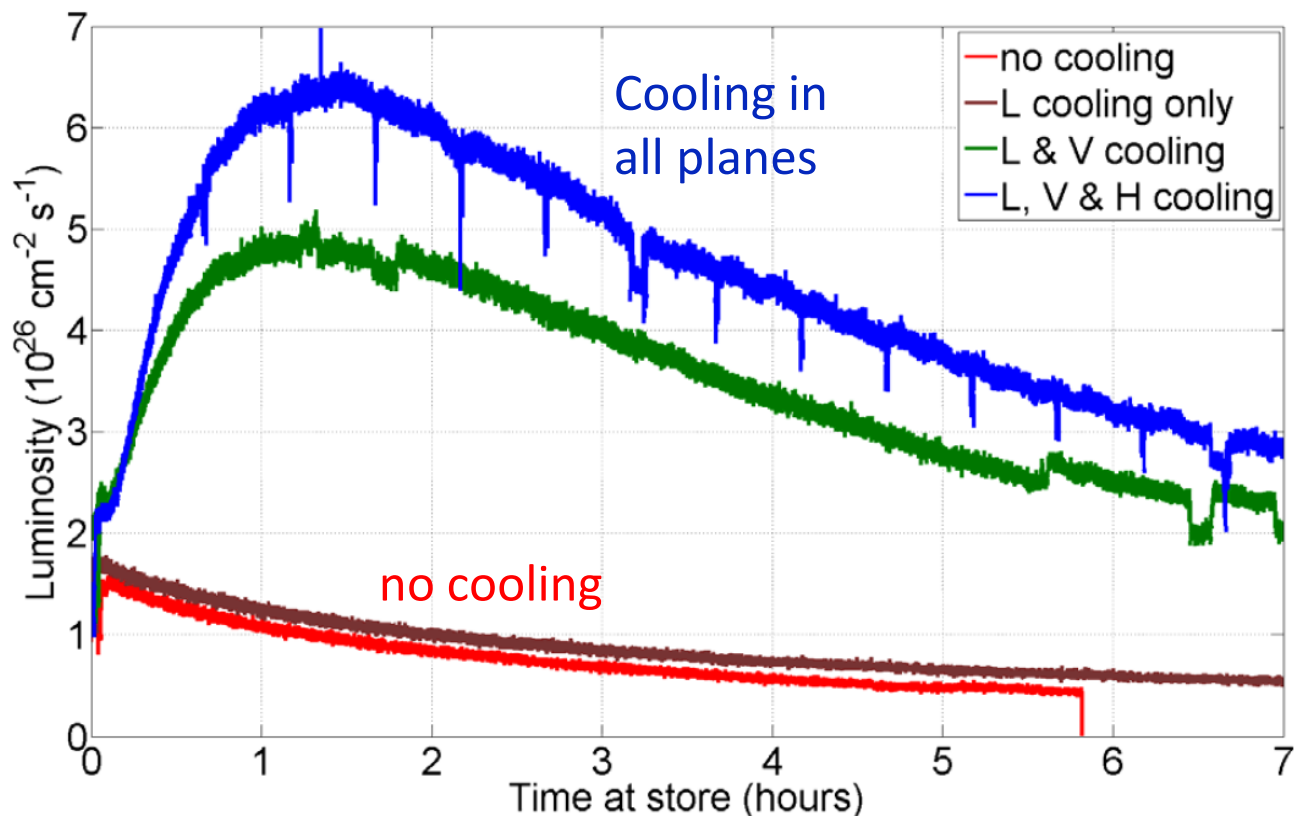
→ treat bunch as part of the beam:  $N \rightarrow \frac{N_b}{B_f} = N_b \frac{C_{LHC}}{4\sigma_z}$

$$\text{Cooling rate per second: } \frac{1}{\tau} = \frac{8W\sigma_z}{N_b C_{LHC}}$$

This simple approximation overestimates the optimum cooling rate by only a factor of 2.

# Stochastic Cooling of Bunched Beams at RHIC

## Luminosity Evolution with and without cooling



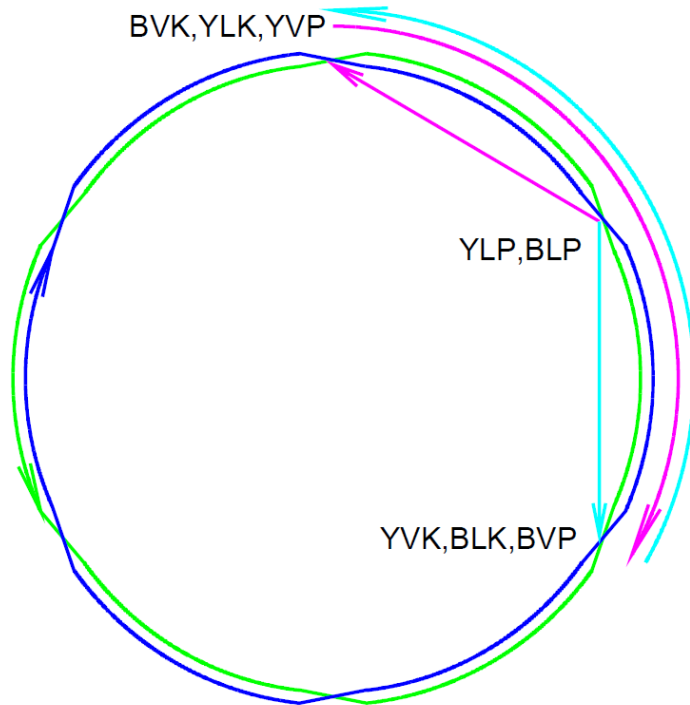
Mike Blaskiewicz

A&T Seminar: <http://indico.cern.ch/conferenceDisplay.py?confId=254917>

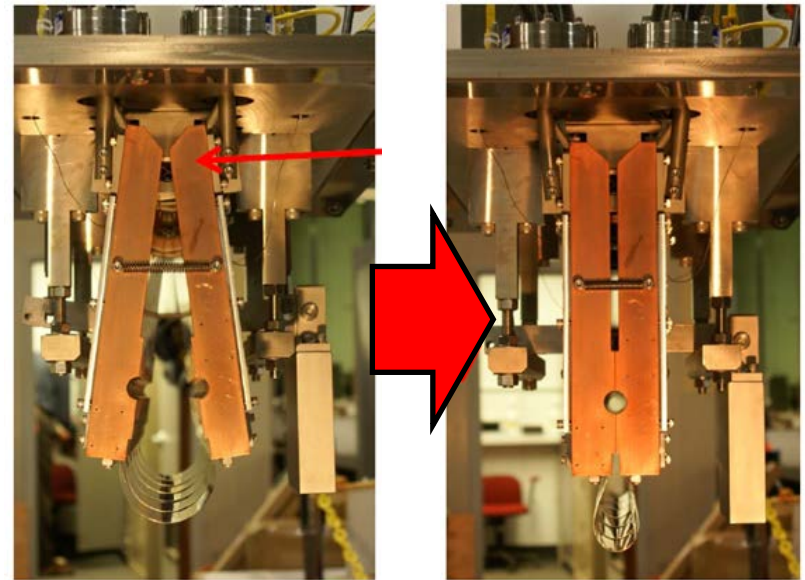
# Stochastic Cooling at RHIC

## Tunnel Layout:

- Transverse signals run backwards in the tunnel.
- Longitudinal signals are sent via microwave link.



Transverse kickers have to be opened at injection due to small aperture of 4.8-7.8GHz cavities. → Impedance problems!

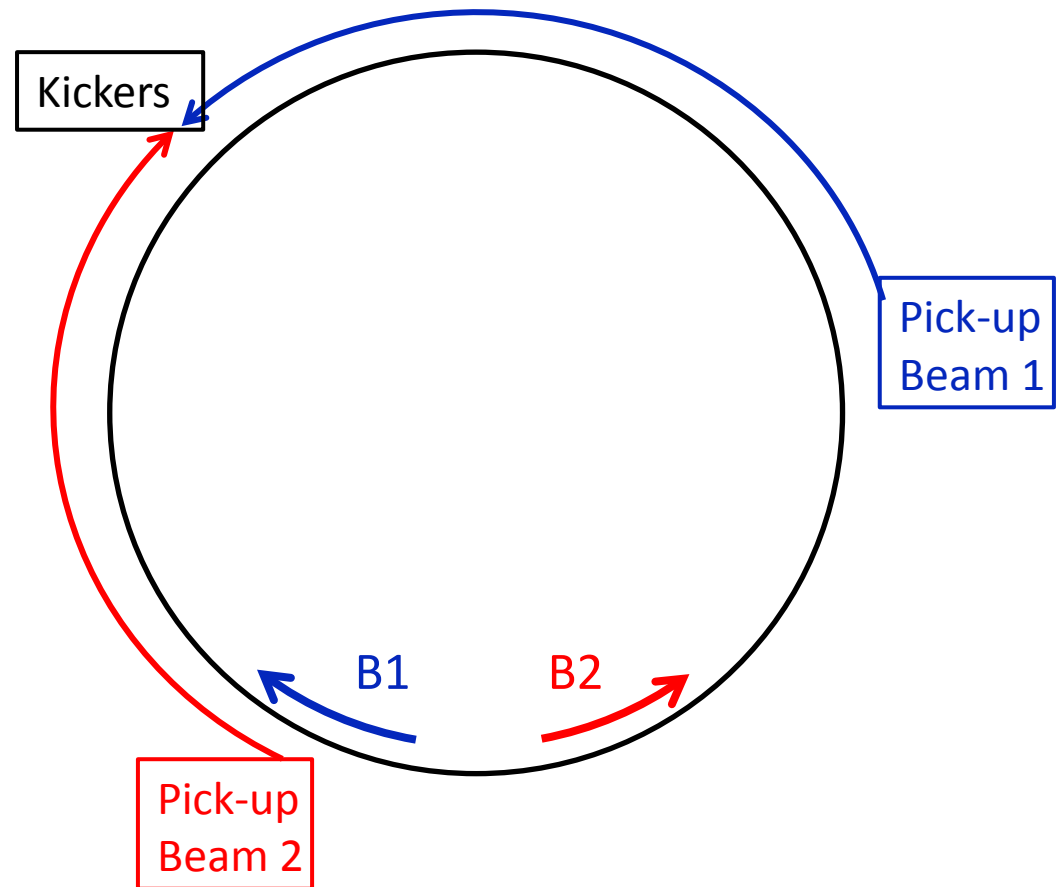


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# Stochastic Cooling System at LHC

- Consider a system similar to the one in RHIC.
- LHC: 27km circumference & 100m underground:
  - New diagonal tunnel for signal too expensive.
  - Microwave-links on surface difficult, due to long distance and weather conditions.
- Signal has to travel backwards in the tunnel.
- Assume 2/3 turn delay between pick-up and kicker.



# Simulations of beam evolution in LHC ring

## Simulations include:

- IBS (various models)
- Burn-off from luminosity production
- Radiation damping and quantum excitation
- Stochastic Cooling

## Simulations require:

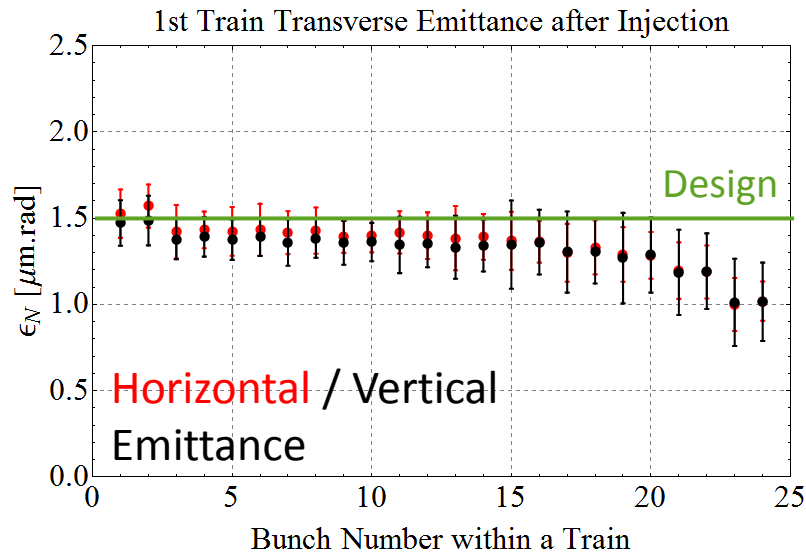
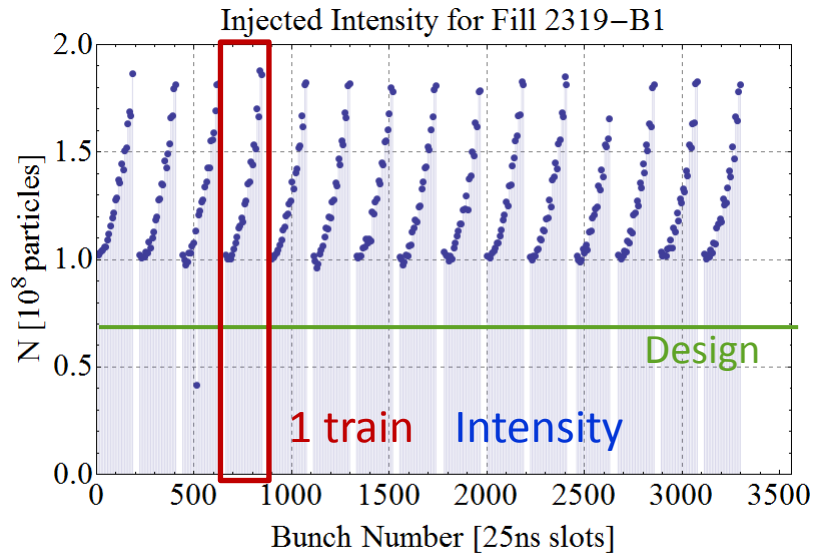
- initial beam parameters (from measurements): e.g. particle type, particles per bunch, emittances, bunch length, RF voltage...
- Properties of stochastic cooling system.

**M. Blaskiewicz's  
Program [1]:  
developed for RHIC**



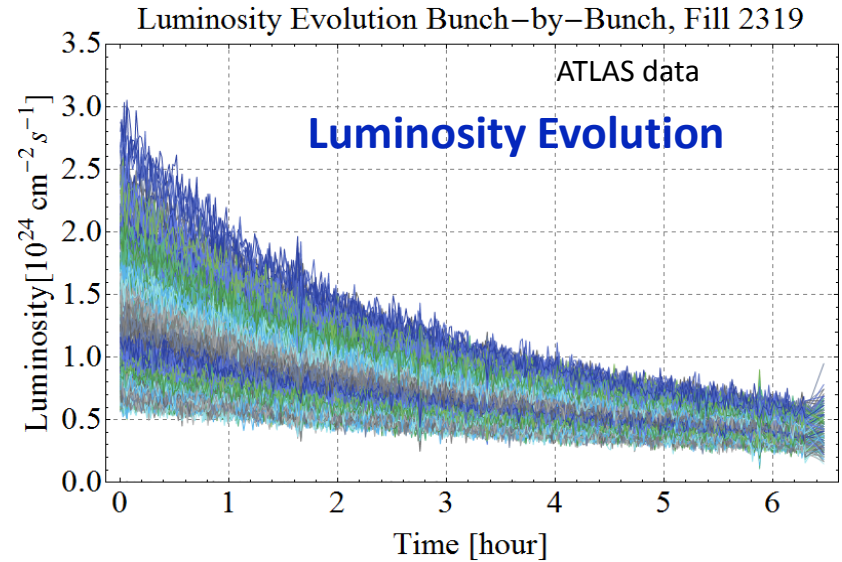
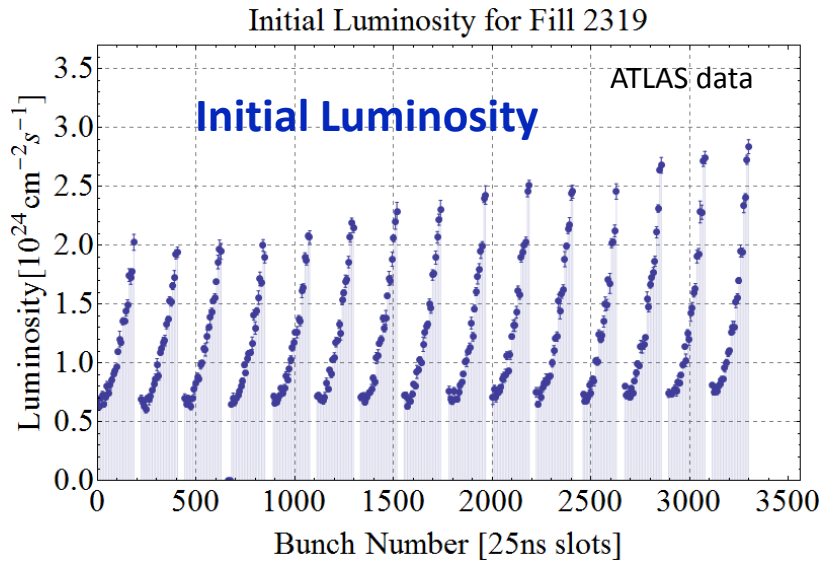
**R. Bruce  
& T. Mertens,  
Collider Time  
Evolution (CTE)  
Program [2]:  
adapted for LHC  
application**

# Bunch-by-Bunch Differences after Injection (450Z GeV)



- Structure within a train (1<sup>st</sup> to last bunch):
  - increase: - intensity  
- bunch length
  - decrease: emittance.
- IBS at the injection plateau of the SPS:
  - while waiting for the 12 injections from the PS to construct a LHC train.
- First injections sit longer at **low energy**
  - strong IBS,
  - emittance growth and particle losses.

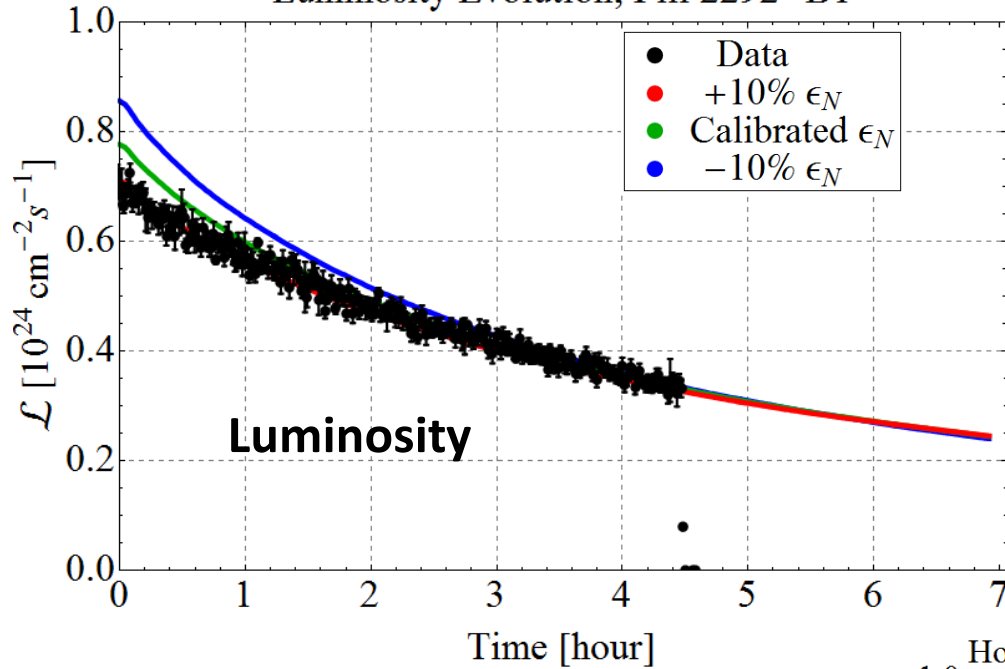
# Luminosity



- Significant bunch-by-bunch structure within a train.
- **Initial values differ by a factor 5-6!**
- Different speed of decay – high initial luminosities decay very fast.

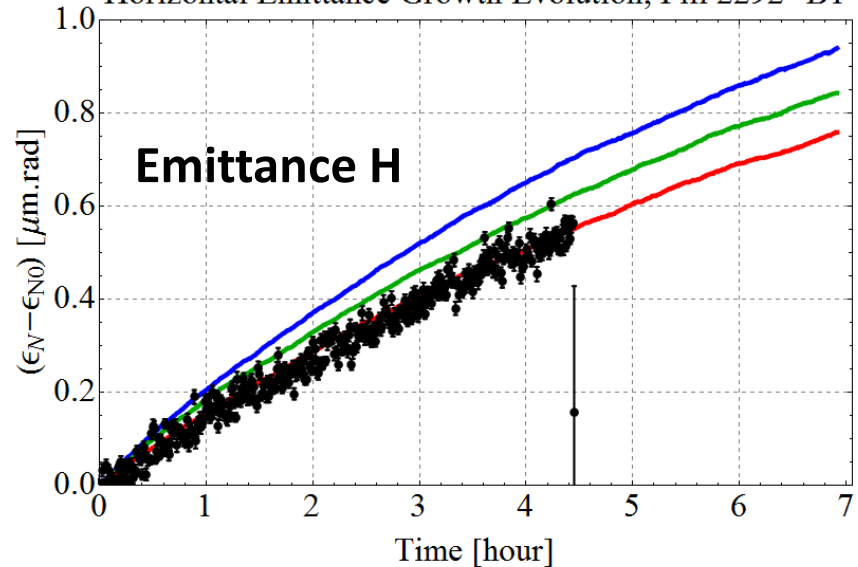
# Evolution in Collisions @ 3.5Z TeV

Luminosity Evolution, Fill 2292-B1



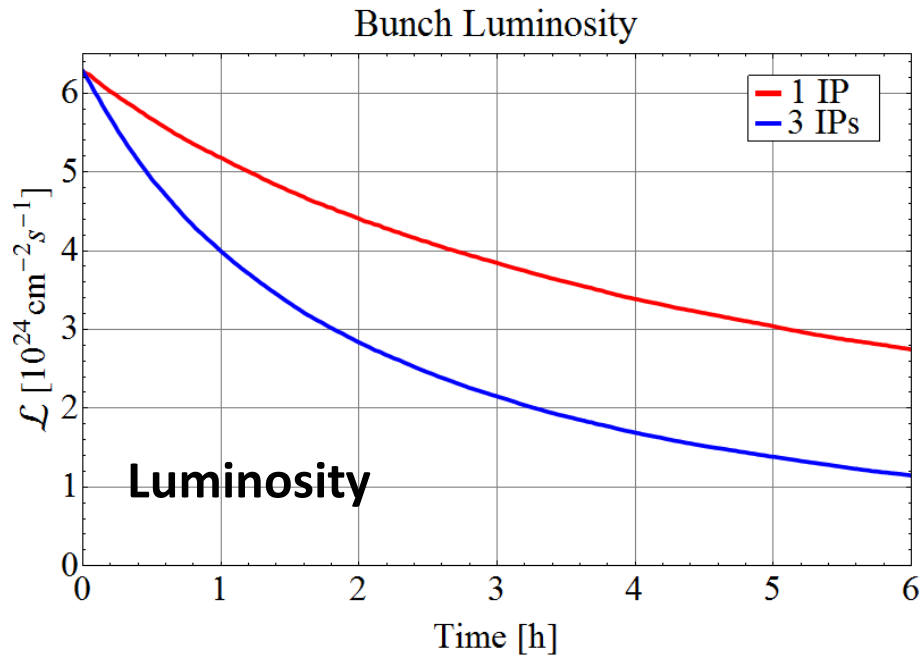
- 2011 Data
- Good agreement between data and simulation.
- Calibration of transverse emittance is difficult.

Horizontal Emittance Growth Evolution, Fill 2292-B1

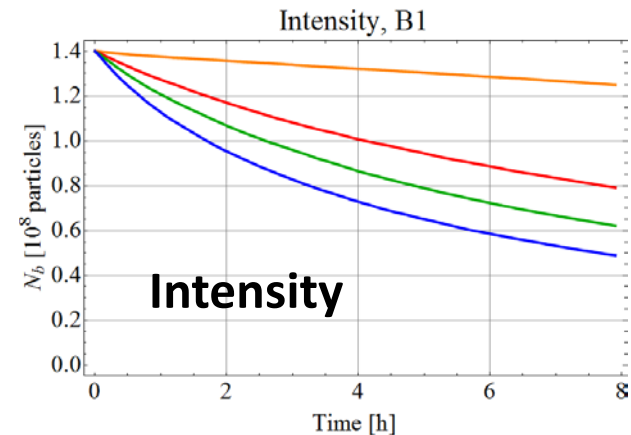
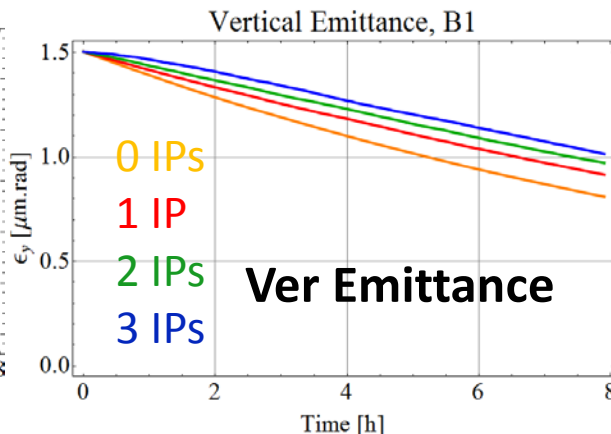
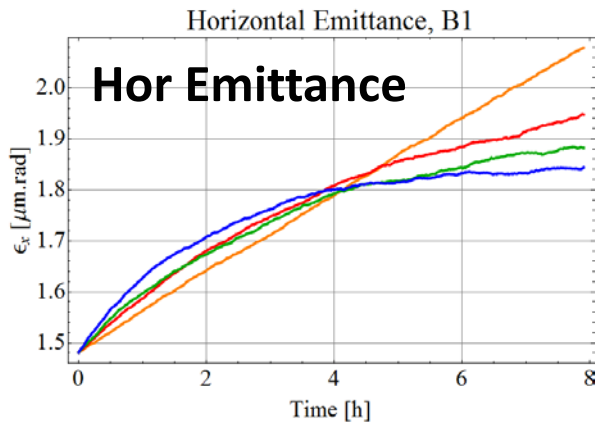




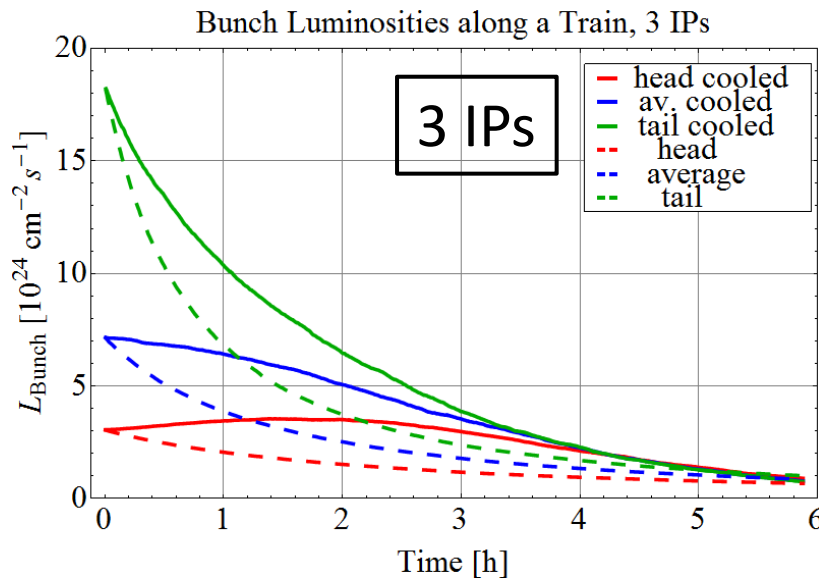
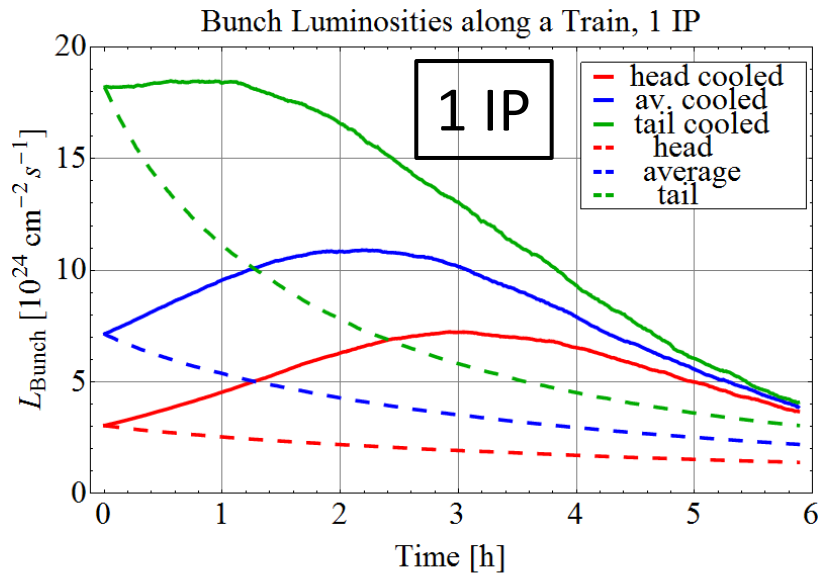
# Potential Beam Evolution @ 7Z TeV



- Simulations [2] with IBS, burn off, radiation damping.
- 3 experiments in collisions lead to very fast burn off:  $\rightarrow$  luminosity  $\frac{1}{2}$ -life  $\approx$  2h.
- Turnaround time  $\approx$  3h.  $\rightarrow$  Longer fills are desired.  $\rightarrow$  Stochastic cooling as possibility to improve fill lifetime.



# Cooling Simulations



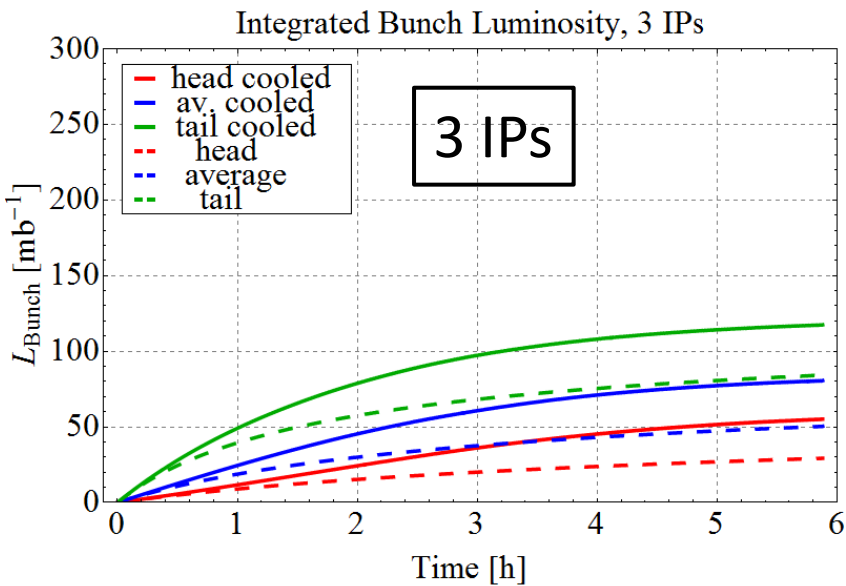
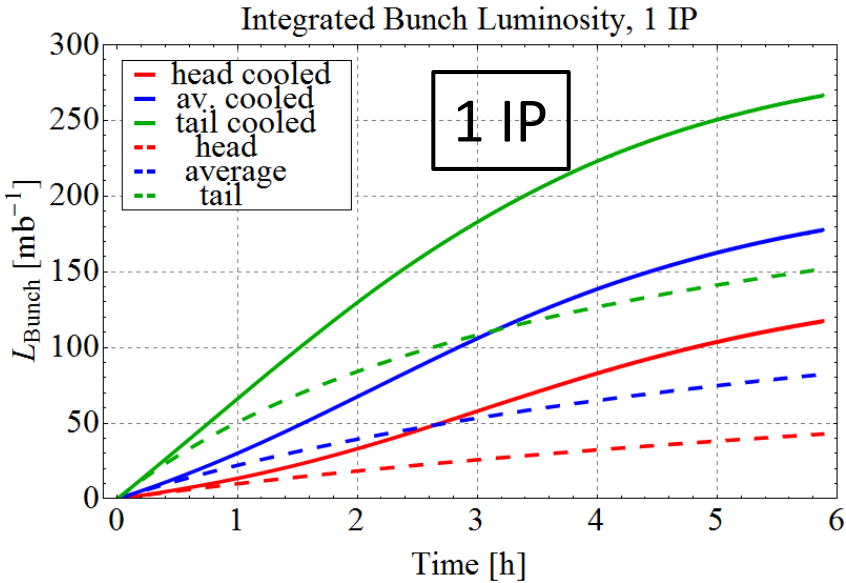
- IBS horizontal growth time  $\approx 8\text{h}$ .
- Radiation damping time  $\approx 13\text{h}$   
 $\rightarrow$  radiation damping not included in the simulations on this slide.
- Assuming a stochastic cooling system with a 5-20GHz bandwidth and average 2013 Pb bunches [4]:

$$T_{\text{cool}} = \frac{N_b C_{\text{LHC}}}{4\sigma_z W} \left[ \frac{M + U}{(1 - \tilde{M}^{-2})^2} \right] \approx 1.8 \text{ h}$$

- First estimate for RMS voltage per cavity (assuming a system with 16 cavities as in RHIC):

$$V_{\text{cavity}} = 2 \text{ kV}$$

# Cooling Simulations



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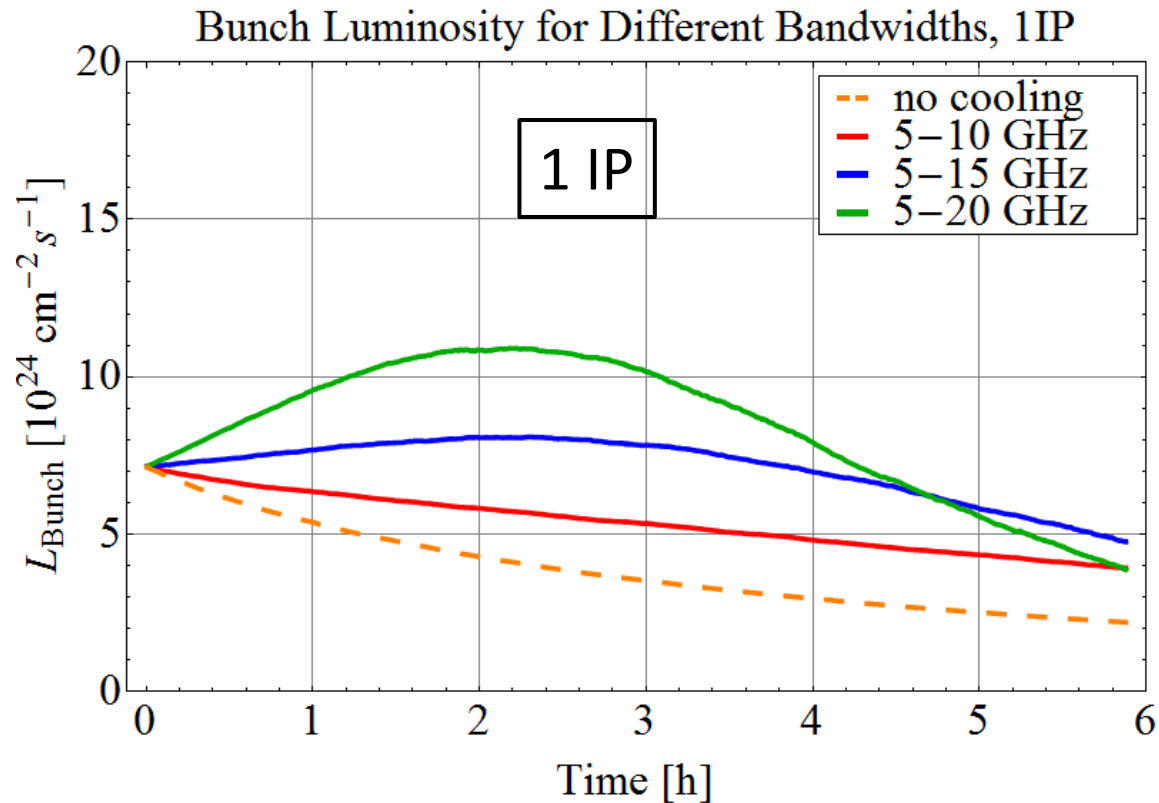
- First estimate for RMS voltage per cavity (assuming a system with 16 cavities as in RHIC):

$$V_{\text{cavity}} = 2 \text{ kV}$$

- Integrated luminosity could be increased by a factor  $\sim 2$ .

# Cooling Simulations

Larger bandwidth and higher upper frequency, lead to higher integrated luminosity.



# Things to be done...

- Find potential locations in the LHC tunnel.
- More detailed simulation and calculations to define required system properties.
- Hardware design challenges to be addressed:
  - Large bandwidth and high upper frequency .
  - Small aperture required → Impedance problems?
  - Compatibility with the proton operation.
- ...

# Conclusions

- **Strong IBS at all energies** leads to emittance growth and particle losses.  
→ Significant bunch-by-bunch differences.
- **Short fills**, due to the high burn off rate with 3 experiments in collisions.
- Stochastic cooling could **equalise bunches and obtain smaller emittances** → higher integrated luminosity.
- First simulation results look promising, studies have just started and are on-going.
  - Challenges in hardware design.

# THANK YOU FOR YOUR ATTENTION

## References:

- [1] M. Blaskiewicz et al., WEM2I05, COOL07 (2007).
- [2] R. Bruce et al., Phys. Rev. ST AB 13, 091001 (2010).
- [3] J. Bjorken, S. Mtingwa, Part. Acc. 13, pp. 115-143 (1983).
- [4] D. Möhl, Lecture Notes in Physics 866, Springer (2013).
- [5] M. Schaumann et al., TUPFI025, IPAC13 (2013).
- [6] M. Schaumann et al., TUPM1HA02, COOL'13 (2013).
- [7] M. Blaskiewicz, CERN A&T Seminar 6/6/2013.

<http://indico.cern.ch/conferenceDisplay.py?confId=254917>

# Design & Current Performance

	Collision (Design)	Injection (2011)	Collision (2011)	Injection (2013)	Collision (2013)
Beam Energy [Z GeV]	7000	450	3500	450	4000
No. Ions per bunch [ $10^8$ ]	0.7	$1.24 \pm 0.30$	$1.20 \pm 0.25$	$1.67 \pm 0.29$	<b><math>1.40 \pm 0.27</math></b>
Transv. normalised emittance [ $\mu\text{m}\cdot\text{rad}$ ]	1.5	---	$1.7 \pm 0.2$	<b><math>1.3 \pm 0.2</math></b>	---
RMS bunch length [cm]	7.94	$8.1 \pm 1.4$	$9.8 \pm 0.7$	$8.9 \pm 0.2$	$9.8 \pm 0.1$
Peak Luminosity [ $10^{27}\text{cm}^{-2}\text{s}^{-1}$ ]	1	---	<b><math>0.4 \pm 0.1</math></b>	---	p-Pb



# Collider Time Evolution (CTE) Program

Processes taken into account:

- COLLISIONS

- user can choose between 2 collision routines:
  - very slow, integrates interaction probability for every particle by sorting particles in opposing beam in discrete bins. **No assumptions on the shape of the beam distribution.**
  - fast routine, **assumes Gaussian transverse distribution** and calculates interaction probability from transverse distribution analytically and uses **global reduction factor** (hourglass and crossing angle) for all particles. **No assumptions on longitudinal distribution.**

- IBS

- rise time calculated using a standard method and modulated to account for non-Gaussian longitudinal profiles
- user can choose between the following methods:
  - Nagaitsev full lattice
  - smooth lattice Piwinski
  - full lattice Piwinski
  - full lattice modified Piwinski
  - full lattice Bane (*not good at injection*)
  - interpolation from tabulated risetimes in external file at given points in emittance-space

- BETATRON MOTION

- SYNCHROTRON MOTION (particles outside RF bucket are lost)

- RADIATION DAMPING and QUANTUM EXCITATION

- transverse aperture cut from COLLIMATION

# Collider Time Evolution (CTE) Program

- Output on a turn-by-turn basis
  - IBS rise times
  - Intensity
  - Transversal and longitudinal emittances
  - Luminosity
- Not Implemented
  - Beam-Beam effects
  - Betatron noise from feedback
    - emittance blow-up
  - RF noise
  - Elastic and inelastic beam gas scattering
    - particle loss and emittance blow-up

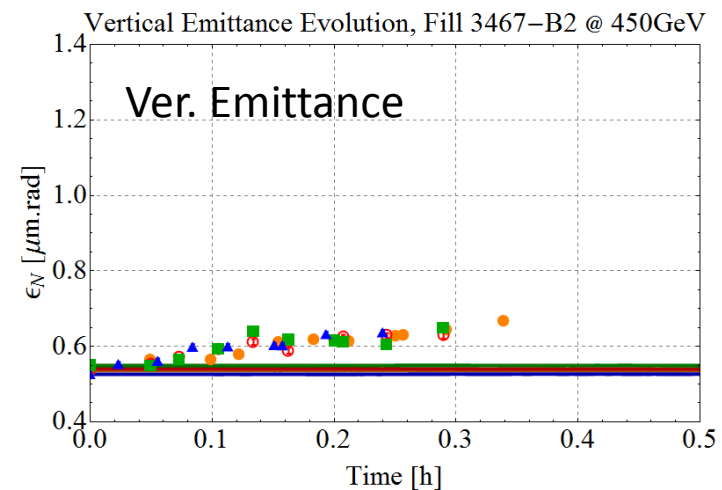
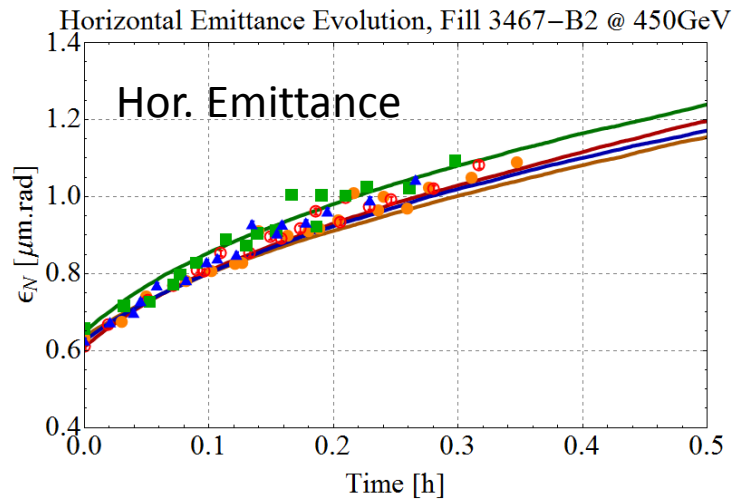
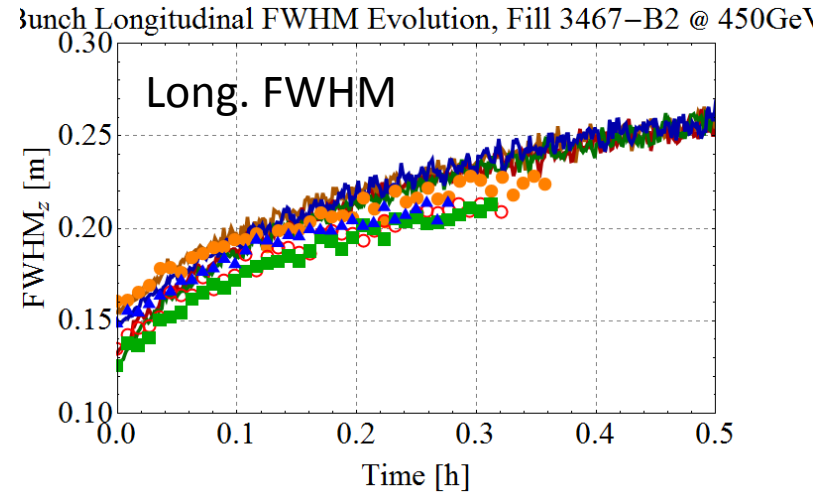
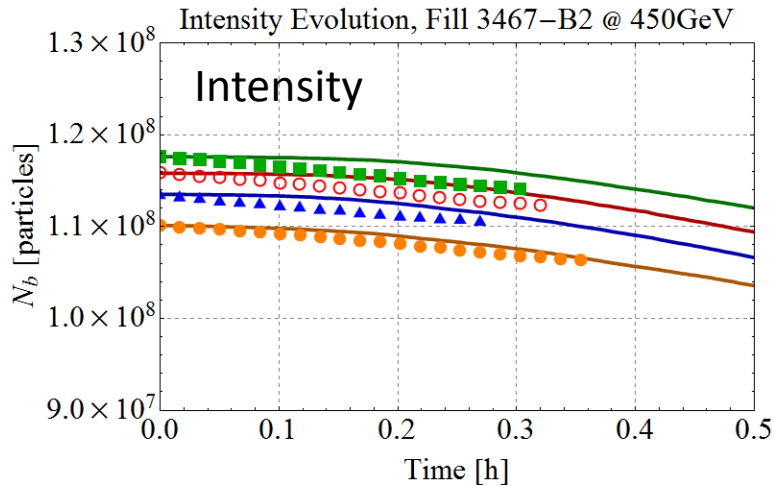
# Beam Evolution at Injection (450Z GeV)

Beams suffer from **strong intra-beam scattering (IBS)**

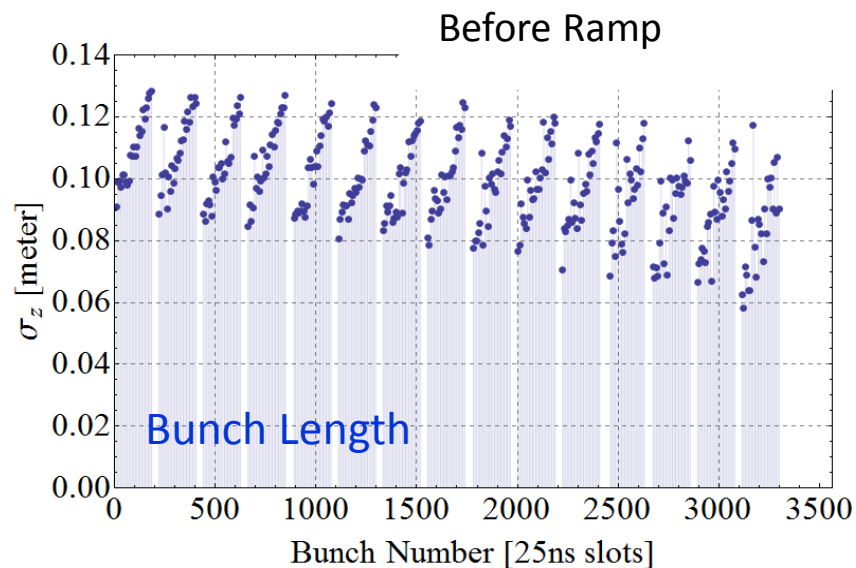
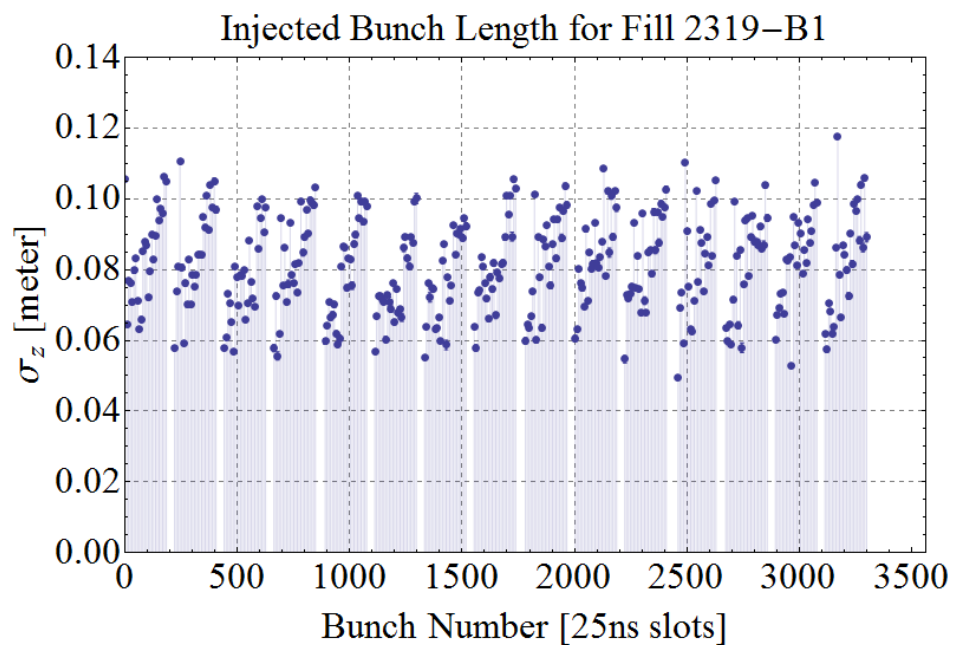
→ **Emittance growth and debunching losses**

Simulations and data are mostly in good agreement.

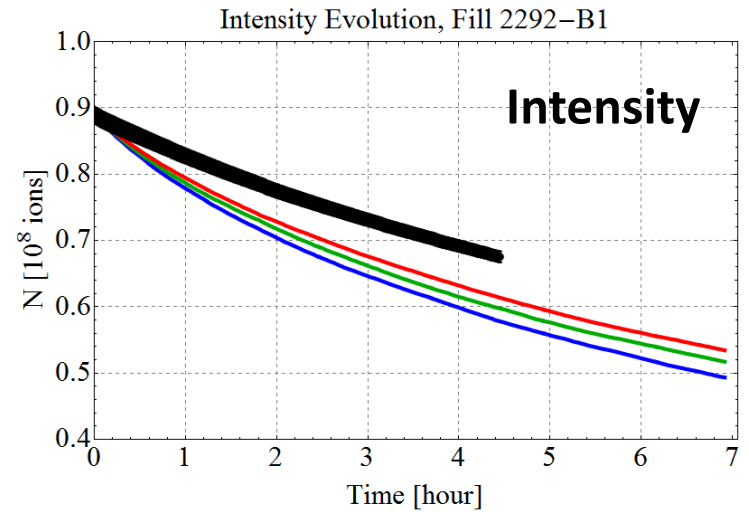
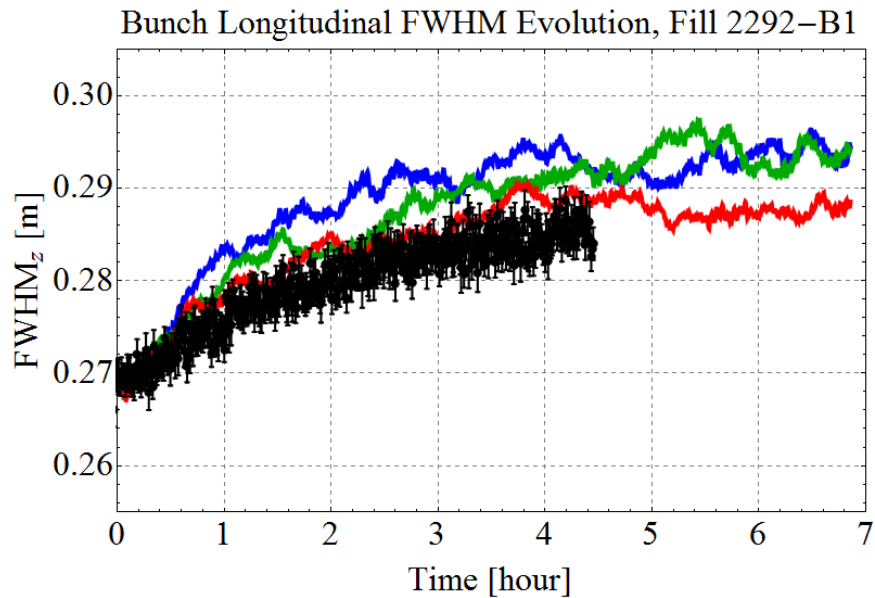
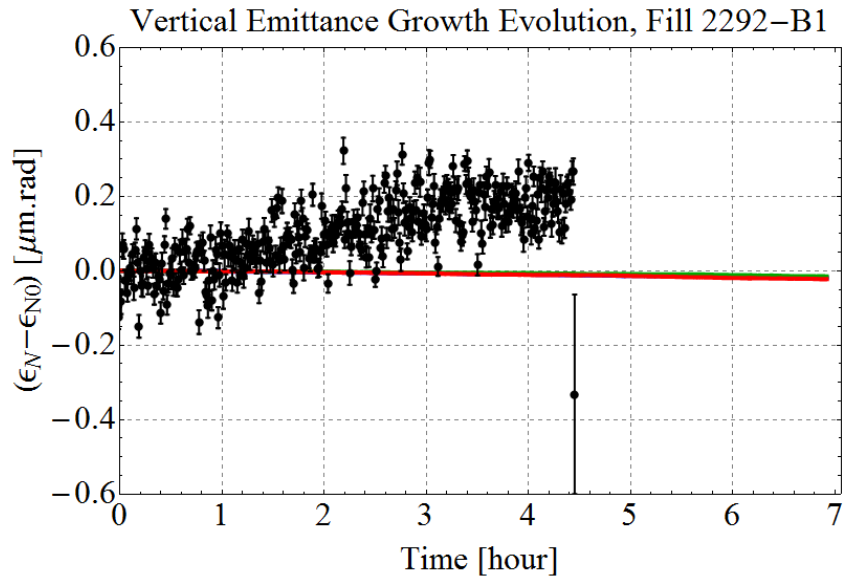
**dots = data**  
**lines = simulation**



# Bunch-by-Bunch Differences after Injection (450Z GeV)



# Evolution in Collisions @ 3.5Z TeV



# Potential Beam Evolution @ 7Z TeV

